

Bond Distance and the Rigid Rotor

Directed Studies Fall 2008

Dale DeWitt

Molecules rotate about a center of mass. The parallel axis theorem

$$I = I_{cm} + Mr^2$$

can be employed, incorporating isotopomers to derive bond lengths. (Mr^2 is indicative of the extra energy required to rotate the body off-center)

There are three elementary bond distance constructs to consider, r_o , r_e , and r_s . Additional constructions can be invoked, but I'll limit this discussion to these basics. r_o , the effective bond distance, is found from ground state moments of inertia. r_e , the equilibrium bond distance, is found at the bottom of the internuclear potential governing vibrational motion in the Born-Oppenheimer approximation. [1] r_s , or substitution, is obtained from isotopomer data sets.

For the simplest case, the rigid rotator, one must procure r_o using a minimum of two vibrational states. [2]

Usually the v_o and v_i vibrational levels are chosen as they are generally the easiest measured.

From p and q branch spacings and rotational quantum mechanics, whereby

$$E_j = BJ(J+1)$$

one can find the required moments of inertia. From the transition v_o to v_i , the $2B$ rotational constant spacing is found. For the R branch

$$\Delta E = \nu_{vib} + 2B(J+1)$$

and for the P branch

$$\Delta E = \nu_{vib} - 2BJ$$

m values, (where $m = J'' + 1$ for the R branch and $m = -J''$ for the P branch) are tabulated with corresponding frequencies, and after fitting a rotational constant B is obtained.

In the ideal case of a non-vibrating (or equilibrium) molecule, bond length is isotopically invariant, being solely determined by charge. [2]

Isotopic substitution will result in a change of the moment of inertia

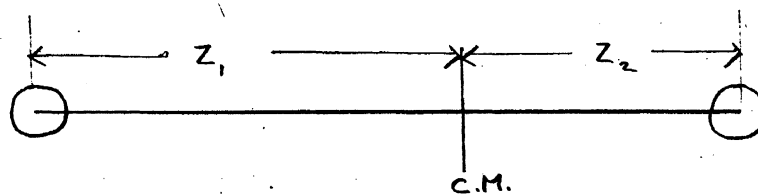
$$\mu' = \frac{M \Delta m}{M + \Delta m}$$

in direct correspondence to a less energetic branch line.

From the center of mass condition

$$m_1 z_1 + m_2 z_2 = 0$$

$$I' - I = (M' - M) \frac{M}{M'} z_2^2 \quad (M = m_1 + m_2)$$

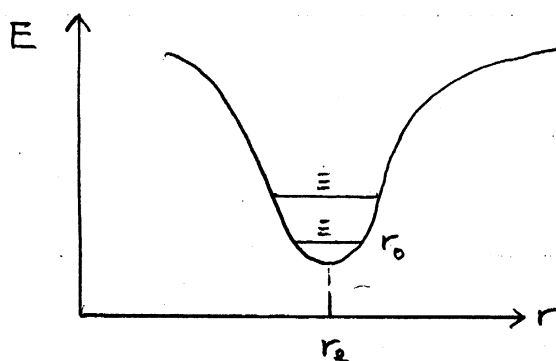


$$|z_2| = \left[\frac{(I' - I)}{M' - M} \frac{M'}{M} \right]^{1/2} \quad \text{or} \quad \left[\frac{\Delta I}{\Delta M} \frac{M'}{M} \right]^{1/2}$$

With the simplest case

$$r_v = \left[\frac{I}{\mu} \right]^{1/2} \quad (\text{where } v = 0, 1, \dots)$$

The r_s method requires one to obtain I_0 values for parent isotopomer and all other substituted daughter species. r_e is independent of isotopic specie, whereas r_0 is inside the potential energy curve and mass dependent. [3]



B_e is an unattainable state but can be approximated from

$$B_v = B_e - \alpha \left(v + \frac{1}{2} \right) \quad (v = 0, 1, \dots)$$

The later term, $\alpha \left(v + \frac{1}{2} \right)$, refers to the vibration-rotation constant. For this r_s method B_0 will be considered isotopically invariant.

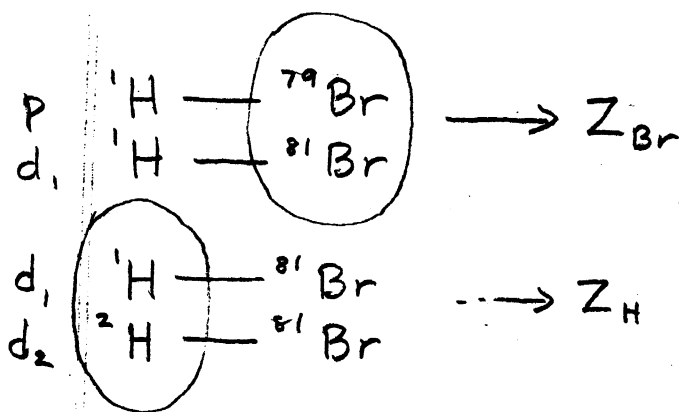
The r_s recipe for a rigid rotor.

Tabulate: ΔI ; ΔM ; M, M' for use in

$$|z_2| = \left[\frac{\Delta I}{\Delta M} \frac{M'}{M} \right]^{1/2} \quad \text{and}$$

calculate p, d_1, d_2 (parent, daughter isotopomers)

For instance



Then $|z_H| + |z_{Br}| = r_s$

Since it is assumed the atomic position do not change under substitution one can literally graft the daughter isotopomer z distances back to the parent.

The r_0 values for HBr

	B_0 (MHz)	I_0
$D^{79}Br$	127357.6334	$6.58932955 \cdot 10^{-47} \text{ Kg m}^2$
$D^{81}Br$	127279.7639	6.593360886.
$H^{79}Br$	250360.12	3.351977211.
$H^{81}Br$	250282.75	3.353013409.

	ΔI
$H^{79}Br - H^{81}Br$	$1.036198 \cdot 10^{-50} \text{ Kg m}^2$
$D^{79}Br - D^{81}Br$	4.031338.
$H^{79}Br - D^{79}Br$	$3.23735234 \cdot 10^{-47} \text{ Kg m}^2$
$H^{81}Br - D^{81}Br$	3.24034748

	ΔM
$H^{79}Br - H^{81}Br$	$3.3177223 \cdot 10^{-27} \text{ Kg}$
$D^{79}Br - D^{81}Br$	3.3177223.
$H^{79}Br - D^{79}Br$	1.6735539.
$H^{81}Br - D^{81}Br$	1.6735539.

	$M \& M'$
$D^{79}Br$	$1.348009286 \cdot 10^{-25} \text{ Kg}$
$D^{81}Br$	1.3812872.
$H^{79}Br$	1.32722186.
$H^{81}Br$	1.36447725.

$$|z_2| = \left[\frac{\Delta I}{\Delta M} \frac{M'}{M} \right]^{1/2}$$

example:

$$\begin{array}{l} p \text{ H}^{79}\text{Br} \\ d_1 \text{ H}^{81}\text{Br} \end{array} \left\{ z_{\text{Br}} \right\} = \left[\frac{1.036198 \cdot 10^{-50} \text{ Kg m}^2}{3.3177223 \cdot 10^{-27} \text{ Kg}} \frac{1.36447725 \cdot 10^{-25} \text{ Kg}}{1.32722186 \cdot 10^{-25} \text{ Kg}} \right]^{1/2}$$

$$|z_{\text{Br}}| = 1.7919 \cdot 10^{-12} \text{ m}$$

$$\begin{array}{l} p \text{ H}^{79}\text{Br} \\ d_2 \text{ D}^{79}\text{Br} \end{array} \left\{ z_{\text{H}} \right\} = \left[\frac{3.23735234 \cdot 10^{-47} \text{ Kg m}^2}{1.6735539 \cdot 10^{-27} \text{ Kg}} \frac{1.348009286}{1.32722186} \right]^{1/2}$$

$$|z_{\text{H}}| = 1.4017 \cdot 10^{-10} \text{ m}$$

	Parent
$r_s = 1.4196 \cdot 10^{-10} \text{ m}$	H ⁷⁹ Br
$= 1.4175 \cdot$	H ⁸¹ Br
$1.4154 \cdot$	D ⁷⁹ Br
$1.4174 \cdot$	D ⁸¹ Br

$$r_0 = \left[\frac{I_0}{\mu} \right]^{1/2}$$

atom
masses

$$H = 1.6653873 \cdot 10^{-27} \text{ kg}$$

$$D = 3.34449401$$

$${}^{79}\text{Br} = 1.314295947 \cdot 10^{-25} \text{ kg}$$

$${}^{81}\text{Br} = 1.343646318 \cdot 10^{-25} \text{ kg}$$

$$\begin{aligned} \mu_{\text{H}^{79}\text{Br}} &= \frac{(1.6653873 \cdot 10^{-27} \text{ kg})(1.314295947 \cdot 10^{-25} \text{ kg})}{(1.6653873 \cdot 10^{-27} \text{ kg} + 1.314295947 \cdot 10^{-25} \text{ kg})} \\ \text{(example)} & \\ &= 1.6653873 \cdot 10^{-27} \text{ kg} \end{aligned}$$

$$r_0 = \frac{I_0}{\mu} = \left[\frac{3.351977211 \cdot 10^{-47} \text{ kg m}^2}{1.6653873 \cdot 10^{-27} \text{ kg}} \right]^{1/2}$$

$$\begin{array}{ll} r_0 = 1.4187 \cdot 10^{-10} \text{ m} & \text{H}^{79}\text{Br} \\ & 1.42769 \cdot 10^{-10} \text{ m} & \text{H}^{81}\text{Br} \\ & 1.42139 \cdot 10^{-10} \text{ m} & \text{D}^{79}\text{Br} \\ & 1.42144 \cdot 10^{-10} \text{ m} & \text{D}^{81}\text{Br} \end{array}$$

Harmony has noted that it is necessary to use two B_v values (i.e., B_0 and B_1) to calculate B_e .

[2, p. 12] For the diatomic

$$B_0 = B_e - \frac{\alpha_e}{2}$$

$$B_1 = B_e - \frac{3\alpha_e}{2}$$

Ignoring the higher order term one finds

$$I_e = \frac{h}{8\pi^2 B_e} = I_0 = \frac{h}{8\pi^2 B_0}$$

For $D^{79}Br$
(example)

$$I_0 = \frac{6.62608 \cdot 10^{-34} \text{ J}\cdot\text{s}}{8\pi^2 (127357.6334 \cdot 10^6 \text{ s}^{-1})} =$$

$$= 6.589341 \cdot 10^{-47} \text{ Kg m}^2$$

$$r_e = \left[\frac{(6.589341 \cdot 10^{-47} \text{ Kg m}^2)}{\left(\frac{(3.344357 \cdot 10^{-27} \text{ Kg})(1.3104705 \cdot 10^{-25} \text{ Kg})}{(3.344357 \cdot 10^{-27} \text{ Kg} + 1.3104705 \cdot 10^{-25} \text{ Kg})} \right)} \right]^{1/2}$$

$$= \left[\frac{6.589341 \cdot 10^{-47} \text{ Kg m}^2}{3.261132 \cdot 10^{-27} \text{ Kg}} \right]^{1/2}$$

$$r_e = 1.42147 \cdot 10^{-10} \text{ m}$$

	(10^{-10}m) $D^{79}\text{Br}$	$D^{81}\text{Br}$
r_0	1.42147	1.42103
r_1	1.435689	1.43568

To estimate r_e from B_0 and B_1 , one must take a mean. For $D^{79}\text{Br}$

$$r_e = \frac{1.42147 + 1.435689}{2} = 1.42858 \cdot 10^{-10} \text{ m}$$

and for $D^{81}\text{Br}$ $r_e = 1.42835 \cdot 10^{-10} \text{ m}$

Tabulating all results

	$H^{79}\text{Br}$	$H^{81}\text{Br}$	$D^{79}\text{Br}$	$D^{81}\text{Br}$
r_0	1.4187	1.42769	1.42139	1.42144
r_s	1.4196	1.4175	1.4154	1.4174
r_e mean			1.4286	1.4284
$r_e(2r_s - r_0)$	1.4205	1.40731	1.40941	1.41336

Solving a simultaneous equation for $B_0 = B_e - \alpha_e$ and $B_1 = B_e = B_e - 3\alpha_e$ gives upon

cancellation of α_e

$$\frac{3}{2} B_0 - \frac{1}{2} B_1 = B_e$$

Plugging in values for $D^{79}\text{Br}$ and $D^{81}\text{Br}$ gives for $D^{79}\text{Br}$

$$B_e = 1.2861299 \cdot 10''$$

$$r_e = \left[\frac{I_e}{\mu} \right]^{1/2} = \left[\frac{6.5250239 \cdot 10^{-47} \text{ Kg m}^2}{3.2614985 \cdot 10^{-27} \text{ Kg}} \right]^{1/2}$$

Then

$$D^{79}\text{Br} = 1.41443 \cdot 10^{-10} \text{ m}$$

$$D^{81}\text{Br} = 1.41484 \cdot 10^{-10} \text{ m}$$

Finding one set of values for α_e for both $D^{79}\text{Br}$ and $D^{81}\text{Br}$ I calculate

$D^{79}\text{Br}$ from B_0

$$127357.6334 \cdot 10^6/s + 1257 \cdot 10^6/s = B_e \quad [4]$$

from B_1 $1.286146334 \cdot 10''$

$$r_e = \left[\frac{6.524940541 \cdot 10^{-47} \text{ Kg m}^2}{3.26149853 \cdot 10^{-27} \text{ Kg}} \right]^{1/2}$$

$$= 1.41442 \cdot 10^{-10} \text{ m}$$

and for $D^{81}\text{Br}$

$$r_e = 1.41446 \cdot 10^{-10} \text{ m}$$

Determining r_e by ignoring the higher term, or vibration-rotation interaction constant, is in many ways avoidance of the problem. Cancellation of the term, given B_0 and B_1 and simultaneously solving for B_e also brings into question just what exactly has been accomplished.

B_e by definition must account for any vibrational disruption of the rotational energies. The best approach would be to solve using an actual higher term which can be extrapolated from R branch compression and P branch rarefaction [6]. Comparing r_e (mean) and $r_e (2r_s-r_0)$, as well as the 'simultaneous equation' method, doesn't shed too much light on the matter.

It is noted the later method, and the r_s method, are smaller valued than the effective method, or r_0 in this case. At least that is logical for deuterium. One would expect a minimum distance at equilibrium as well as inconsistent data with the hydrogen isotopomer.

References

1. H. W. Kroto, *Molecular Rotation Spectra*, John Wiley & Sons, London 1975, p. 156
2. Marlin D. Harmony, *Molecular Structure Determination From Spectroscopic Data Using Scaled Moments of Inertia, Vibrational Spectra and Structure*, (J. R. Durig, Ed.) Vol. 24, Amsterdam: Elsevier, 1999.
3. J. Michael Hollas, *Modern Spectroscopy*, John Wiley & Sons, New York 1996, p. 116.
4. C. H. Townes: A. L. Schawlow, *Microwave Spectroscopy*, McGraw-Hill, New York 1955, p. 639.
5. J. Michael Hollas, *Basic Atomic and Molecular Spectroscopy*, Wiley-Interscience 2002, p. 130.
6. David P. Shoemaker, et al., *Experiments in Physical Chemistry* 5th Ed., McGraw-Hill, New York 1989, p. 463 Experiment 38 Vibrational-Rotational Spectra of HCl and DCl.
7. Marlin D. Harmony, *Introduction to Molecular Energies and Spectra*, Holt, Rinehart and Winston New York, 1972.